

THE COMPLEX MONGE-AMPÈRE EQUATION AND AN APPLICATION TO UNIFORMISATION OF SURFACES

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ABSTRACT. We prove that a complete noncompact Kähler surface with positive and bounded sectional curvature is biholomorphic to $\mathbb{C}^2$. This result confirms a special case of Yau’s conjecture that a complete noncompact Kähler n -manifold with positive holomorphic bisectional curvature is biholomorphic to $\mathbb{C}^n$. In contrast to all known results on Yau’s conjecture, we do not need additional assumptions on the global/asymptotic geometry of the Kähler surface apart from completeness. Towards this end, we prove that the integral of the square of the Ricci form of a complete Kähler surface with positive sectional curvature is finite. The work of Chen and Zhu shows that this latter result implies that the surface is biholomorphic to $\mathbb{C}^2$. The main new idea is the construction of a Lipschitz continuous plurisubharmonic weight function with finite Monge-Ampère mass. This weight function is obtained by solving a complex Monge-Ampère equation.

1. INTRODUCTION

Let (X, ω) be a complete Kähler manifold of complex dimension n and positive holomorphic bisectional curvature (which we denote by $BK_\omega > 0$). If X is compact, it is biholomorphic to $\mathbb{P}^n$ by the resolution of the Frankel conjecture by Siu-Yau [35] and Mori [30]. If X is noncompact, a longstanding conjecture of Yau predicts that X must be biholomorphic to $\mathbb{C}^n$ [38]. The corresponding statement under the stronger assumption of positive sectional curvature was previously conjectured by Green and Wu [13].

Yau’s conjecture has been settled under additional hypotheses, typically involving volume growth and curvature decay / finiteness of certain curvature integrals [12, 27, 28, 24, 25, 26, 3]. In [21], Liu proved that if X is as above and has maximal volume growth, then it is indeed biholomorphic to $\mathbb{C}^n$. Subsequently Lee and Tam [18], building on earlier work by Chau, Tam and others (cf. [3] and references therein), proved the same result using the Kähler-Ricci flow. Interestingly, the other extreme case of *minimal* volume growth has also been settled, at least in complex dimension 2: In [5], Chen-Zhu proved that if X is a complete noncompact Kähler manifold with $BK > 0$ and $p \in X$, then there exists $C > 0$ such that

$$Vol(B(p, r)) \geq Cr^n,$$

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where $n = \dim_{\mathbb{C}} X$. In a later work [8], they proved that if X has positive *sectional* curvature and minimal volume growth, then it is biholomorphic to an affine algebraic variety. If $\dim_{\mathbb{C}} X = 2$, it follows that X is biholomorphic to $\mathbb{C}^2$ by a classical result of Ramanujam [33]. In a different direction, Chen-Zhu [6] proved that if X has positive and *bounded* sectional curvature, and satisfies

$$\int_X \text{Rc}_{\omega}^n < \infty,$$

then X is biholomorphic to a quasi-projective variety (cf. [26, 36]). Again, if $\dim_{\mathbb{C}} X = 2$, M is biholomorphic to $\mathbb{C}^2$.

Yet another conjecture of Yau [39], refined by Yang [37], states that if (X^n, ω) has non-negative bisectional curvature and $o \in X$ is a given point, then there exists a constant $C > 0$ such that for any $r > 0$ and $k = 1, \dots, n$,

$$r^{2k-2n} \int_{B(o,r)} \text{Rc}_{\omega}^k \wedge \omega^{n-k} \leq C.$$

These are higher-dimensional versions of the classical Cohn-Vossen inequality [9]. If X^n has positive sectional curvature, then the $k = 1$ case of the conjecture follows from the following a priori estimate of Petrunin [32] and rescaling: There exists a dimensional constant $C(n)$ such that for any $o \in X$, the scalar curvature S_{ω} satisfies

$$\int_{B(o,1)} S_{\omega} \omega^n < C.$$

Using our weight function and the method of [6] we can establish the remaining case of Yau's Cohn-Vossen type conjecture for Kähler surfaces with positive sectional curvature.

Theorem 1. *Let (X, ω) be a complete noncompact Kähler surface with positive sectional curvature. Then*

$$\int_X \text{Rc}_{\omega}^2 < \infty.$$

Combining Theorem 1 with the result of Chen and Zhu mentioned above [6] we obtain the following special case of the conjectures of Green-Wu and Yau.

Theorem 2. *Let (X, ω) be a complete noncompact Kähler surface with positive and bounded sectional curvature. Then X is biholomorphic to $\mathbb{C}^2$.*

Note that the positive sectional curvature hypothesis implies that X is diffeomorphic to $\mathbb{R}^4$, by the Gromoll-Meyer theorem [14]. It also implies that X is a Stein manifold equipped with a smooth strictly plurisubharmonic (psh) exhaustion function that is uniformly Lipschitz, by the work of Green-Wu [12]. Neither of these facts is known to hold under the weaker assumption $BK_{\omega} > 0$.

The key new technical input in this paper, which may be of independent interest, is the following construction of a uniformly Lipschitz psh weight function with finite Monge-Ampère mass.

Theorem 3. *Let (X^n, ω) be a complete, non-compact n -dimensional Kähler manifold with a smooth exhaustion function $\rho : X \rightarrow \mathbb{R}$ with uniformly bounded gradient. Suppose $BK_\omega > 0$. Then there exists a uniformly Lipschitz, strictly plurisubharmonic function ϕ on X such that*

$$\int_X (\sqrt{-1} \partial \bar{\partial} \phi)^n < \infty.$$

In particular, in view of the theorem of Green and Wu [12], such a ϕ exists on any complete, non-compact Kähler manifold with positive sectional curvature.

Note that there is no dimension restriction in the above theorem.

The present work originated from an attempt to use non-smooth weights to resolve Yau's uniformisation conjecture for Kähler surfaces with positive sectional curvature, without imposing an upper bound. Although that objective remains unfulfilled, a finite stratification result for such manifolds will be presented in a forthcoming paper now in preparation.

2. CONSTRUCTION OF A WEIGHT FUNCTION WITH BOUNDED MASS

We prove Theorem 3 in this section. Without loss of generality, we may also assume that $\rho \geq 0$, and that $|\nabla \rho|_\omega \leq 1$. We fix a point $o \in \rho^{-1}(0)$. Note that in the case of positive sectional curvature, by Green and Wu's work one can choose ρ to be strictly convex, and hence o be the unique point in $\rho^{-1}(0)$. Next, we set

$$B_R := \{x \in M \mid \rho(x) < R\},$$

and $S_R := \partial B_R$. Since ρ is a strictly psh exhaustion function, B_R is a strictly pseudoconvex compact subset of X . Let $R_\nu \rightarrow \infty$ be a sequence of regular values of ρ . Set $B_\nu = B_{R_\nu}$ and $S_\nu = \partial B_\nu$. The main idea is to solve a complex Monge-Ampère equation on B_ν with a rapidly decaying right hand side and take a limit of the solutions. The key point is to obtain apriori gradient estimates so as to obtain a Lipschitz function in the limit. We first need to construct a rapidly decaying function at infinity that is smaller in some precise sense than the curvature decay at infinity. For any point $p \in M$ we let

$$\Lambda(x) := \inf_{u, v \in T_x M} \frac{\text{Rm}(u, v, v, u) + \text{Rm}(u, Jv, Jv, u)}{|u \wedge v|^2}.$$

By our hypothesis, $\Lambda(p) > 0$ for all p . We then have the following elementary observation.

Lemma 4. *There exists a smooth exhaustion function $F : [0, \infty) \rightarrow \mathbb{R}$ and a constant $C > 0$ such that the following properties hold:*

(1)

$$\int_X e^{-F(\rho)/n} \omega^n < \infty,$$

(2)

$$e^{-F(\rho)} \omega^n < \frac{1}{2} (\sqrt{-1} \partial \bar{\partial} \rho)^n,$$

and

(3) For every $x \in X$,

$$\|\nabla F(\rho)\|_\omega e^{-F(\rho)/n}(x) \leq \frac{n}{4} \Lambda(x)$$

Proof. Fix a point $o \in X$. Let $g : [0, \infty) \rightarrow (0, \infty)$ be a smooth strictly decreasing function satisfying

$$g(r) < \left(\inf_{x \in S_r} \min \left(e^{-d_\omega(o, x)}, \frac{(\sqrt{-1} \partial \bar{\partial} \rho)^n}{2\omega^n}(x) \right) \right)^{1/n}.$$

Note that by the exhaustive nature of ρ , S_r is non-empty for each $r \geq 0$. Moreover, the right-hand-side is some continuous, strictly positive function in r , and hence this is always possible. One can easily see that $\lim_{r \rightarrow \infty} g(r) = 0$.

We now define

$$F(\rho) = -n \ln \left(c - \int_0^\rho G(s) ds \right),$$

where

$$c = \int_0^\infty G(s) ds,$$

and we choose a smooth, strictly positive function G so that

- $G(s) < -g'(s)$ for all s , and
- $G(r) < \frac{1}{4} \inf_{x \in S_r} \Lambda(x)$.

From the first condition it follows that

$$\int_r^\infty G(s) ds < g(r).$$

In particular G is integrable with the total integral $c < g(0)$. It is now easy to check that $F(\rho)$ satisfies all the required properties. $\square$

Next, by [15], there exists a strictly psh $u_\nu \in C^\infty(B_\nu) \cap C^0(\overline{B_\nu})$ such that

$$\begin{cases} (\sqrt{-1} \partial \bar{\partial} u_\nu)^n = e^{-F(\rho)} \omega^n \\ u_\nu|_{S_\nu} = \rho = R_\nu. \end{cases}$$

By the comparison principle and our choice of F we have that $u_\nu \geq \rho$. The key point is the following:

Lemma 5. For all ν ,

$$\sup_{B_\nu} |\nabla u_\nu|_\omega^2 \leq 1.$$

Proof. We first use the maximum principle (cf. Blocki [2]) to reduce the estimate to the boundary. For ease of notation, we will drop the subscript ν . We also let $\omega' := \sqrt{-1}\partial\bar{\partial}u$ and denote all quantities (such as the Laplacian) associated to ω' with a prime. Suppose $x_0 \in A_R$ is a maxima for $|\nabla u|^2$. We choose normal coordinates for ω so that $\sqrt{-1}\partial\bar{\partial}u$ is diagonal with entries $(\lambda_1, \dots, \lambda_n)$. We also denote by h the endomorphism of $T^{1,0}X$ given by $h_k^i = g^{i\bar{j}}u_{k\bar{j}}$. We now compute at p ,

$$\begin{aligned}\Delta'|\nabla u|^2 &= u^{i\bar{j}}\partial_i\partial_{\bar{j}}(g^{k\bar{l}}u_ku_{\bar{l}}) \\ &= u^{i\bar{j}}R_{i\bar{j}k\bar{l}}g^{k\bar{q}}g^{p\bar{l}}u_ku_{\bar{l}} + \text{Tr}(h) + u^{i\bar{j}}g^{k\bar{l}}u_{ik}u_{\bar{j}\bar{l}} + u^{i\bar{j}}g^{k\bar{l}}u_{k\bar{j};i}u_{\bar{l}} + u^{i\bar{j}}g^{k\bar{l}}u_ku_{\bar{l};i}.\end{aligned}$$

Note that the final two terms involve the third derivatives of u . One can use the equation to simplify these terms. Indeed, taking log and differentiating the equation with respect to ∂_k we obtain

$$u^{i\bar{j}}u_{i\bar{j};k} = -\partial_k F(\rho) + g^{p\bar{q}}g_{p\bar{q};k} = -\partial_k F(\rho)$$

since we are working with normal coordinates for ω at p . We can also similarly write a formula for the $\bar{l}$ -derivative. Note also that $u_{i\bar{j};k} = u_{k\bar{j};i}$ and so we get that

$$\Delta'|\nabla u|^2 = u^{i\bar{j}}R_{i\bar{j}k\bar{l}}g^{k\bar{q}}g^{p\bar{l}}u_ku_{\bar{l}} + \text{Tr}(h) + u^{i\bar{j}}g^{k\bar{l}}u_{ik}u_{\bar{j}\bar{l}} - 2\langle \nabla u, \nabla F(\rho) \rangle_\omega.$$

In normal coordinates the third term takes the forms

$$u^{i\bar{j}}g^{k\bar{l}}u_{ik}u_{\bar{j}\bar{l}} = \sum_{i,k} \frac{|u_{ik}|^2}{\lambda_i} \geq 0.$$

So finally we have

$$\begin{aligned}\Delta'|\nabla u|^2 &\geq u^{i\bar{j}}R_{i\bar{j}k\bar{l}}g^{k\bar{q}}g^{p\bar{l}}u_ku_{\bar{l}} - 2|\nabla F(\rho)||\nabla u| \\ &\geq \Lambda|\nabla u|^2\text{Tr}(h^{-1}) - 2|\nabla F(\rho)||\nabla u| \\ &\geq \left(n\Lambda e^{F(\rho)/n} - 2\frac{F'(\rho)}{|\nabla u|}\right)|\nabla u|^2,\end{aligned}$$

where we used the arithmetic-geometric mean inequality in the final line. Without loss of generality we may assume that $|\nabla u|^2(p) \geq 1$. Plugging this in we obtain

$$\Delta'|\nabla u|^2 \geq (n\Lambda e^{F(\rho)/n} - 2F'(\rho))|\nabla u|^2 > 0,$$

which contradicts the maximum principle. To summarise, we have proven that

$$\sup_{\overline{B_R}} |\nabla u|^2 = \sup_{\partial B_R} |\nabla u|^2.$$

Clearly we only need to bound the normal derivative at the boundary. In the neighbourhood of the boundary, note that by the comparison principle,

$$\rho \leq u_\nu \leq R_\nu.$$

The normal derivative $\nabla_n u = \langle \nabla u, n \rangle$, where $n = -\nabla \rho / |\nabla \rho|$ is the inward pointing normal, then clearly satisfies

$$\nabla_n \rho \leq \nabla_n u \leq 0,$$

and hence we have the required bound. $\square$

Proof of Theorem 3. We let $\phi_\nu(x) = u_\nu(x) - u_\nu(o)$. Then $|\nabla \phi_\nu| \leq 1$. Moreover, since $\phi_\nu(o) = 0$, for on any compact set K , there exists a constant C_K such that $\|\phi_\nu\|_{C^1(K)} \leq C_K$. By a standard Arzela-Ascoli and diagonal argument, after passing to a subsequence, ϕ_ν uniformly converge on compact sets to a Lipschitz function ϕ with $\text{Lip}(\phi) \leq 1$. By the continuity of the Monge-Ampère operator under uniform limits (cf. [11, pg. 147]), ϕ solves the Monge-Ampère equation

$$(\sqrt{-1}\partial\bar{\partial}\phi)^n = e^{-F(\rho)}\omega^n,$$

and hence by construction,

$$\int_M (\sqrt{-1}\partial\bar{\partial}\phi)^n < \infty.$$

$\square$

3. PROOF OF THEOREM 1

3.1. Smoothing by heat flow. Recall that since $\text{Rc}_\omega \geq 0$, there exists a unique, positive, symmetric and stochastically complete heat kernel $H(x, y, t)$. We let

$$u(x, t) = \int_X H(x, y, t) \phi(y) dy.$$

Then $u(x, t)$ is a solution to the heat equation with initial condition $u(x, 0) = \phi(x)$. By [31], $u(x, t)$ is strictly psh for each t . We let $\psi(x) := u(x, 1)$. Then by [5], there exists a constant $A > 0$ such that we have the following estimates:

$$|\nabla u|, t|\sqrt{-1}\partial\bar{\partial}u| \leq A.$$

We will work with both the non-smooth weight ϕ and the smooth weights $u(x, t)$. To switch back and forth between the two weights, we need the following crucial estimate. The argument appears to be standard but we found it through Chatgpt 5.0.

Lemma 6. *There exists a dimensional constant $c(n)$ such that for any $0 \leq t_1 \leq t_2$,*

$$u(x, t_2) \leq u(x, t_1) + Ac\sqrt{t_2 - t_1},$$

where A is the Lipschitz constant for ϕ . In particular, $\psi(x) \leq \phi(x) + cA$.

Proof. We have the representation formula

$$u(x, t_2) = \int_X H(x, y, t_2 - t_1) u(y, t_1) dy.$$

From the fact that $u(x, t_1)$ is also Lipschitz with the Lipschitz constant A , and stochastic completeness, it follows that

$$u(x, t_2) \leq u(x, t_1) + A \int_X H(x, y, t) d(x, y) dy,$$

where we set $t = t_2 - t_1$. So it suffices to estimate the integral on the right. By the fundamental Li-Yau gradient estimates [20], we have the following Gaussian estimate:

$$H(x, y, t) \leq \frac{C}{|B(x, \sqrt{t})|} e^{-c \frac{d^2(x, y)}{t}},$$

for some $c < 1/4$. Now for integers $k \geq 0$, consider the annuli

$$A_k = B(x, (k+1)\sqrt{t}) \setminus B(x, k\sqrt{t}).$$

Then

$$\begin{aligned} \int_X H(x, y, t) d(x, y) dy &\leq \sum_k \int_{A_k} H(x, y, t) d(x, y) dy \\ &\leq \frac{C\sqrt{t}}{|B(x, \sqrt{t})|} \sum_k (k+1) e^{-ck^2} |A_k|. \end{aligned}$$

But now by the Bishop-Gromov inequality,

$$\frac{|A_k|}{|B(x, \sqrt{t})|} \leq \frac{|B(x, (k+1)\sqrt{t})|}{|B(x, \sqrt{t})|} \leq \omega_4 (k+1)^4,$$

and so

$$\int_X H(x, y, t) d(x, y) dy \leq C\omega_2 \sqrt{t} \sum_{k=1}^{\infty} (k+1)^5 e^{-ck^2} \leq c\sqrt{t}.$$

□

Remark 7. An interesting question is whether ψ continues to have finite Monge-Ampère mass.

We also need the following basic observation on constructing suitable cut-off functions.

Lemma 8. *Let (X, ω) satisfy $BK_\omega \geq 0$. Fix $o \in X$. Then there exist $0 < \theta < 1$, $A > 0$ and $a_0 > 0$ such that the following holds: for all $a > a_0$ there exist a smooth function $\chi_a : X \rightarrow [0, 1]$ having the following properties:*

- (1) $\chi_a \equiv 1$ on $B(o, \theta a)$ and $\text{Supp}(\chi_a) \subset B(o, \theta^{-1}a)$.
- (2) There exists a constant A such that

$$|\nabla \chi_a|, |\sqrt{-1} \partial \bar{\partial} \chi_a| \leq \frac{A}{a}.$$

Proof. This is standard, so we only sketch the proof. Let $u(x, t)$ solve the heat equation with $u(x, 0) = d(x, o)$, and let $\eta(x) = u(x, 1)$. Then there exists a constant $C > 1$ such that

- $C^{-1}(1 + d(x, o)) \leq \eta(x) \leq C(1 + d(x, o))$.
- $|\nabla \eta|, |\sqrt{-1}\partial\bar{\partial}\eta| < C$.

Now let $\chi : \mathbb{R} \rightarrow [0, 1]$ be the usual cut-off function such that $\chi \equiv 1$ on $t \leq 1$ and $\text{Supp}(\chi) \subset (-\infty, 2]$. Then

$$\chi_a(x) := \chi\left(\frac{\eta(x)}{a}\right)$$

does the job with $\theta = (2C)^{-1}$ and $a_0 = 2C$. $\square$

3.2. Estimates on holomorphic sections. We first recall the following classical theorem of Hormander and Andreotti-Vesentini.

Theorem 9 (Hormander, Andreotti-Vesentini). *Let (X, ω) be a complete Kähler manifold, let u be a smooth function on X , and let L be a holomorphic line bundle equipped with a smooth hermitian metric h such that the curvature satisfies*

$$\sqrt{-1}\partial\bar{\partial}u + \sqrt{-1}\Theta_h + \text{Ric}(\omega) \geq c(x)\omega,$$

for some continuous function $c : X \rightarrow (0, \infty)$. Suppose we have an L -valued $(0, 1)$ form β satisfying $\bar{\partial}\beta = 0$ and

$$\int_X \frac{\|\beta\|^2}{c} e^{-u} \omega^n < \infty.$$

Then there exists a unique $\xi \in \Gamma(L)$ satisfying $\bar{\partial}\xi = \beta$ and the L^2 -estimate

$$\int_X \frac{|\xi|^2}{c} e^{-u} \omega^n \leq \int_X \frac{\|\beta\|^2}{c} e^{-u} \omega^n.$$

With ϕ and ψ as in the previous section, we consider Hermitian metrics $h_{q\phi} = e^{-q\phi}(\omega^n)^{-1}$ and $h_{q\psi} = e^{-q\psi}(\omega^n)^{-1}$ on the canonical bundle K_M . We denote the norms simply as $\|\cdot\|_{q\phi}$ and $\|\cdot\|_{q\psi}$ respectively (suppressing in particular the dependence on ω)

Lemma 10. *Let u be a solution to the heat equation with initial data $u(x, 0) = \phi$ as above. Then there exists a $q \gg 1$, a non-trivial holomorphic section of $s \in H^0(X, K_X)$ and a constant $C > 0$ such that for all $0 \leq t \leq 1$,*

$$\int_X \|s\|_{qu}^2 \omega^2, \int_X \|\nabla_{qu}s\|_{qu}^2 \omega^2 < C.$$

Proof. As above, let $\psi = u(x, 1)$. The proof of existence of holomorphic sections, L^2 -integrable with respect to the weight ψ is standard. Nevertheless, we include an outline for the convenience of the reader. Fix a point $o \in X$. By scaling we may assume that there exist holomorphic coordinates (z^1, z^2) on the ball $B(o, 2)$. Let χ be a cut-off function with support in $B(o, 2)$ such that $\chi \equiv 1$ on $B(o, 1)$. Let $\sigma = \chi \cdot dz^1 \wedge dz^2$ and $\beta = \bar{\partial}_{K_M} \sigma$, where $\bar{\partial}_{K_M}$ in the $\bar{\partial}$ operator on the canonical line bundle. Then β is a $\bar{\partial}$ -closed $(1, 0)$ K_M -valued form. We now apply the Hörmander-Andreotti-Vesentini

Theorem 9 above with the Hermitian metric $h = (\omega^2)^{-1}$ and the weight function

$$\tilde{\psi} = q\psi + 4\chi \cdot \log |z|^2.$$

For $q \gg 1$, clearly $\sqrt{-1}\partial\bar{\partial}\tilde{\psi} > c\omega$, for some continuous function $c \in C^0(X)$. Without loss of generality we can assume that $c < 1$ and that $c > \delta$ on U . Moreover

$$\sqrt{-1}\partial\bar{\partial}\tilde{\psi} + \Theta_h + \text{Rc}(\omega) = \sqrt{-1}\partial\bar{\partial}\tilde{\psi} > 0,$$

and so there exists a solution $\xi \in \Gamma(K_X)$ to $\bar{\partial}_{K_X}\xi = -\beta$ satisfying

$$\int_X |\xi|^2 e^{-q\psi - 4\chi \log |z|^2} \omega^2 < C.$$

In particular, $\xi(o) = 0$. Clearly, we also have that

$$\int_X |\xi|^2 e^{-q\psi} \omega^2 < \infty.$$

Finally let $s = \sigma + \xi$. Then $s \in H^0(M, K_X)$ and satisfies

$$\int_X |s|^2 e^{-q\psi} \omega^2 < C,$$

for some $C > 0$. By the estimate in Lemma 6, we see that there exists a dimensional constant A such that for all $0 \leq t \leq 1$,

$$\psi(x) \leq u(x, t) + A\sqrt{1-t}.$$

Then

$$\int_X \|s\|_{qu}^2 \omega^2 = \int_X |s|^2 e^{-qu} \omega^2 \leq e^{Aq} C.$$

To obtain a gradient bound, we make use of the following Böchner-Weitzenböck formula:

$$\Delta \|s\|_{q\psi}^2 = \|\nabla_{q\psi} s\|_{q\psi}^2 - q\Delta\psi + S_\omega \|s\|_{q\psi}^2 \geq \|\nabla_{q\psi} s\|_{q\psi}^2 - A\|s\|_{q\psi}^2,$$

where S_ω is the scalar curvature of ω . Let χ_a be the family of cut-off functions from Lemma 8. Then multiplying by χ_a^2 and integrating by parts we obtain

$$\begin{aligned} \int_X \chi_a^2 \|\nabla_{q\psi} s\|_{q\psi}^2 \omega^n &\leq A \int_X \chi_a^2 \|s\|_{q\psi}^2 \omega^n + \int_X \Delta \chi_a^2 \|s\|_{q\psi}^2 \omega^n \\ &\leq C'. \end{aligned}$$

This implies an upper bound for $\int_X \|\nabla_{qu} s\|_{q\phi}^2 \omega^2$ for all $t \in [0, 1]$: Since

$$\nabla_{qu} s = \nabla_{q\psi} s + q s d(u - \psi),$$

and u, ψ are Lipschitz with Lipschitz constant A , we get

$$|\nabla_{qu} s|_{qu}^2 \leq 2|\nabla_{q\psi} s|_{qu}^2 + 2q^2 A^2 |s|_{qu}^2.$$

Lemma 6 then once again gives

$$\int_X \|\nabla_{qu} s\|_{qu}^2 \omega^n \leq C'' \left(\int_X \|\nabla_{q\psi} s\|_{q\psi}^2 \omega^n + \int_X \|s\|_{q\psi}^n \omega^2 \right) \leq C'' C,$$

for some C'' depending only on q and A .

□

As a consequence we obtain the following estimate:

Lemma 11. *Let $s \in H^0(X, K_X)$ be the holomorphic section constructed in Lemma 10 above. Then there exists a constant C such that for all $0 \leq t \leq 1$,*

$$\int_X \|s\|_{qu}^2 \sqrt{-1} \partial \bar{\partial} u \wedge \omega \leq C.$$

Proof. Let χ_a be the family of cut-off functions from Lemma 8, and let

$$I_t(a) := \int_X \chi_a \|s\|_{qu}^2 \sqrt{-1} \partial \bar{\partial} u \wedge \omega.$$

The required estimate follows from the following claim: There exists a constant $C > 0$ such that $I_t(a) \leq C$ for all $0 \leq t \leq 1$ and for all $a \geq a_0$, where a_0 is as in Lemma 8. First assume that $t > 0$. Integrating by parts we see that

$$I_t(a) = \int_X \|s\|_{qu}^2 \partial \chi_a \wedge \bar{\partial} u \wedge \omega + \int_X \chi_a \partial \|s\|_{qu}^2 \wedge \bar{\partial} u \wedge \omega.$$

The first term is clearly uniformly bounded by Lemma 10 and the uniform Lipschitz control on u . For the second term we note that if α, β are two $(1, 0)$ forms then $|\alpha \wedge \bar{\beta} \wedge \omega| \leq |\alpha| |\beta| \omega^2$, and since u is Lipschitz with Lipschitz constant bounded by A , we then have that

$$\left| \int_X \chi_a \partial \|s\|_{qu}^2 \wedge \bar{\partial} u \wedge \omega \right| \leq A \int_X \chi_a |\nabla \|s\|_{qu}^2| \omega^2.$$

By the Kato inequality and the AM-GM inequality, for any section ξ of a vector bundle,

$$(1) \quad |\nabla |\xi|^2| \leq 2|\xi| |\nabla \xi| \leq |\xi|^2 + |\nabla \xi|^2.$$

Applying this to $\xi = s$ in the above expression we see that

$$\int_X \chi_a |\nabla \|s\|_{qu}^2| \omega^2 \leq \int_X \chi_a \|s\|_{qu}^2 \omega^2 + \int_X \chi_a \|\nabla_{qu} s\|_{qu}^2 \omega^2.$$

Each term is uniformly bounded by the Lemma above. Finally, the claimed estimate at $t = 0$ follows from standard Bedford-Taylor theory (cf. [11]) and the fact that $u(x, t)$ converges uniformly to $\phi(x)$ on compact sets.

□

3.3. Proof of Theorem 1. Let $s \in H^0(X, K_X)$ as in Lemma 10. For $\varepsilon > 0$ we consider the $(1, 1)$ current

$$\zeta_\varepsilon(s) := \sqrt{-1} \partial \bar{\partial} \log(\|s\|_{q\phi}^2 + \varepsilon^2) + q \sqrt{-1} \partial \bar{\partial} \phi.$$

Lemma 12. $\zeta_\varepsilon(s)$ is a closed, positive $(1, 1)$ current satisfying

$$\zeta_\varepsilon(s) \geq \frac{\|s\|_{q\phi}^2}{\|s\|_{q\phi}^2 + \varepsilon^2} \text{Ric}(\omega).$$

In particular, by Fatou's lemma,

$$\int_X \text{Ric}(\omega)^2 \leq \liminf_{\varepsilon \rightarrow 0^+} \int_X \zeta_\varepsilon(s) \wedge \zeta_\varepsilon(s),$$

where the wedge product on the right is interpreted in the Bedford-Taylor sense.

Proof. Let χ be a strongly positive $(1, 1)$ form on X with compact support. Without loss of generality, we may assume that χ is supported on a coordinate neighbourhood U . Let $v_\varepsilon = \log(\|s\|_{q\phi}^2 + \varepsilon^2)$. By Bedford-Taylor theory,,

$$\int_X \zeta_\varepsilon(s) \wedge \chi = \int_X (v_\varepsilon + q\phi) \sqrt{-1} \partial \bar{\partial} \chi.$$

Let ϕ_δ be a smoothening of ϕ (say via convolution with an approximation to identity) such on the support of χ we have that $\phi_\delta \rightarrow \phi$ uniformly and in $W^{1,p}$ for all $p > 1$. Note also that ϕ_δ is pluri-subharmonic for all $\delta > 0$. Let $v_{\varepsilon,\delta} := \log(\|s\|_{q\phi_\delta}^2 + \varepsilon^2)$, and

$$\zeta_{\varepsilon,\delta}(s) := \sqrt{-1} \partial \bar{\partial} v_{\varepsilon,\delta} + q \sqrt{-1} \partial \bar{\partial} \phi_\delta$$

in U . By the dominated convergence theorem and integration by parts,

$$\int_X (v_\varepsilon + q\phi) \sqrt{-1} \partial \bar{\partial} \chi = \lim_{\delta \rightarrow 0} \int_X \chi \wedge \zeta_{\varepsilon,\delta}(s).$$

Next, for any smooth function f and any constant c , an elementary computation shows that on the set where $f \neq 0$,

$$\sqrt{-1} \partial \bar{\partial} \log(f + c) \geq \frac{f}{f + c} \sqrt{-1} \partial \bar{\partial} \log f.$$

Taking $f = \|s\|_{q\phi_\delta}^2$ and $c = \varepsilon^2$ we get that on $X \setminus \{s = 0\}$,

$$\begin{aligned} \zeta_{\varepsilon,\delta}(s) &\geq \frac{\|s\|_{q\phi_\delta}^2}{\|s\|_{q\phi_\delta}^2 + \varepsilon^2} \sqrt{-1} \partial \bar{\partial} \log \|s\|_{q\phi_\delta}^2 + q \sqrt{-1} \partial \bar{\partial} \phi_\delta \\ &= \frac{\|s\|_{q\phi_\delta}^2}{\|s\|_{q\phi_\delta}^2 + \varepsilon^2} \text{Ric}(\omega) + \frac{\varepsilon^2 q}{\|s\|_{q\phi_\delta}^2 + \varepsilon^2} \sqrt{-1} \partial \bar{\partial} \phi_\delta \\ &\geq \frac{\|s\|_{q\phi_\delta}^2}{\|s\|_{q\phi_\delta}^2 + \varepsilon^2} \text{Ric}(\omega). \end{aligned}$$

But then clearly the inequality also holds on all of X . Integrating against χ , and letting $\delta \rightarrow 0$ we obtain the required lower bound for $\zeta_\varepsilon(s)$. The second part follows from the Bedford-Taylor theory. Indeed if S, T are closed positive $(1, 1)$ currents with bounded local potentials such that $T \geq S$, then $T^2 \geq S^2$, where the wedge product is interpreted in the Bedford-Taylor sense. □

Theorem 1 is now a consequence of the following:

Lemma 13. *For each $k = 1, 2$,*

$$\int_X \zeta_\varepsilon(s)^k \wedge (\sqrt{-1}\partial\bar{\partial}\phi)^{2-k} \leq q^k \int_X (\sqrt{-1}\partial\bar{\partial}\phi)^2.$$

Proof. We first prove the Lemma for $k = 1$. Let χ_a be the family of cut-off functions from before (cf. Lemma 8). It is enough to prove that

$$I(a) := \int_X \chi_a^3 \sqrt{-1}\partial\bar{\partial} \ln(\|s\|_{q\phi}^2 + \varepsilon^2) \wedge \sqrt{-1}\partial\bar{\partial}\phi \xrightarrow{a \rightarrow \infty} 0.$$

From the definition of wedge products in the Bedford-Taylor theory, we can integrate by parts, and obtain

$$\begin{aligned} I(a) &= \int_X \log\left(1 + \frac{\|s\|_{q\phi}^2}{\varepsilon^2}\right) \sqrt{-1}\partial\bar{\partial}\chi_a^3 \wedge \sqrt{-1}\partial\bar{\partial}\phi \\ &\leq \frac{C}{a} \int_X \log\left(1 + \frac{\|s\|_{q\phi}^2}{\varepsilon^2}\right) \omega \wedge \sqrt{-1}\partial\bar{\partial}\phi \\ &\leq \frac{C}{a\varepsilon^2} \int_X \|s\|_{q\phi}^2 \omega \wedge \sqrt{-1}\partial\bar{\partial}\phi \\ &\leq \frac{C'}{a\varepsilon^2} \end{aligned}$$

for some constant C' independent of a by Lemma 11. Note that we used the estimates on χ_a from Lemma 8 in the second line, and the elementary inequality $\log(1+x) \leq x$ for $x \geq 0$ in the third line.

Next, we consider $k = 2$. It is enough to prove that

$$J(a) := \int_X \chi_a^3 \sqrt{-1}\partial\bar{\partial} \log(\|s\|_{q\phi}^2 + \varepsilon^2) \wedge \zeta_\varepsilon(s) \xrightarrow{a \rightarrow \infty} 0.$$

Once again by Bedford-Taylor theory, $J(a) = \lim_{t \rightarrow 0^+} J_t(a)$, where

$$J_t(a) = \int_X \chi_a^3 \sqrt{-1}\partial\bar{\partial} \log(\|s\|_{q\phi}^2 + \varepsilon^2) \wedge \zeta_{\varepsilon,t}(s), \text{ and}$$

$$\zeta_{\varepsilon,t}(s) := \sqrt{-1}\partial\bar{\partial} \log(\|s\|_{qu}^2 + \varepsilon^2) + q\sqrt{-1}\partial\bar{\partial}u.$$

It suffices to prove that there exists a constant C (independent of t and a) such that $|J_t(a)| \leq Ca^{-1}$. Integrating by parts as before and using the upper bound on $\sqrt{-1}\partial\bar{\partial}\chi_a$ and $|\nabla\chi_a|$, we have

$$\begin{aligned} |J_t(a)| &\leq \frac{C}{a} \int_X \chi_a \log\left(1 + \frac{\|s\|_{qu}^2}{\varepsilon^2}\right) \zeta_{\varepsilon,t}(s) \wedge \omega \\ &= \frac{C}{a} \int_X \chi_a \log\left(1 + \frac{\|s\|_{qu}^2}{\varepsilon^2}\right) \sqrt{-1}\partial\bar{\partial} \log\left(1 + \frac{\|s\|_{qu}^2}{\varepsilon^2}\right) \wedge \omega \\ &\quad + \frac{Cq}{a} \int_X \chi_a \log\left(1 + \frac{\|s\|_{qu}^2}{\varepsilon^2}\right) \sqrt{-1}\partial\bar{\partial}u \wedge \omega. \end{aligned}$$

The second integral is uniformly in t and a by Lemma 11 and the elementary inequality $\ln(1+x) \leq x$ for $x \geq 0$. For the first integral, we note that it

is non-negative and so it is enough to obtain a uniform (in t and a) upper bound. Integrating by parts,

$$\begin{aligned} \text{first integral} &= - \int_X \chi_a \sqrt{-1} \partial \log \left(1 + \frac{\|s\|_{qu}^2}{\varepsilon^2} \right) \wedge \bar{\partial} \log \left(1 + \frac{\|s\|_{qu}^2}{\varepsilon^2} \right) \wedge \omega \\ &\quad - \int_X \log \left(1 + \frac{\|s\|_{qu}^2}{\varepsilon^2} \right) \sqrt{-1} \partial \chi_a \wedge \bar{\partial} \log \left(1 + \frac{\|s\|_{qu}^2}{\varepsilon^2} \right) \wedge \omega \\ &\leq \frac{C}{a} \int_X \log \left(1 + \frac{\|s\|_{qu}^2}{\varepsilon^2} \right) \left| \nabla \log \left(1 + \frac{\|s\|_{qu}^2}{\varepsilon^2} \right) \right| \omega^2, \end{aligned}$$

where we used the fact that the first integral in the expression above is non-negative and once again the elementary observation that $|\alpha \wedge \beta \wedge \omega| \leq |\alpha| |\beta| \omega^2$ for $(1, 0)$ forms α, β along with the estimate $|\nabla \chi_a| \leq C a^{-1}$. Finally,

$$\begin{aligned} \int_X \log \left(1 + \frac{\|s\|_{qu}^2}{\varepsilon^2} \right) \left| \nabla \log \left(1 + \frac{\|s\|_{qu}^2}{\varepsilon^2} \right) \right| \omega^2 &\leq \frac{1}{\varepsilon^2} \int_X \frac{\|s\|_{qu}^2}{\|s\|_{qu}^2 + \varepsilon^2} |\nabla \|s\|_{qu}^2| \omega^2 \\ &\leq \frac{1}{\varepsilon^2} \int_X |\nabla \|s\|_{qu}^2| \omega^2 \\ &\leq \frac{1}{\varepsilon^2} \int_X \|s\|_{qu} \|\nabla_{qu} s\|_{qu} \omega^2 \\ &\leq \frac{1}{\varepsilon^2} \left(\int_X \|s\|_{qu}^2 \omega^2 \right)^{1/2} \left(\int_X \|\nabla_{qu} s\|_{qu}^2 \omega^2 \right)^{1/2} \\ &\leq C \end{aligned}$$

for some uniform constant C , independent of a and t by Lemma 10. Note that we once again used the Kato inequality (1) in the third line. $\square$

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