

Spectral Mixture Modeling with Laboratory Near-Infrared Data I: Insights into Compositional Analysis of Europa

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Abstract

Europa’s surface composition and physical characteristics are commonly constrained using spectral deconvolution through linear mixture (LM) modeling and radiative transfer-based (RT) intimate mixture modeling. Here, I compared the results of these two spectral modeling— LM versus RT— against laboratory spectra of water (H₂O) ice and sulfuric acid octahydrate (SAO; H₂SO₄·8H₂O) mixtures measured at near-infrared wavelengths (~ 1.2 – $2.5\ \mu\text{m}$) with grain sizes of 90–106 μm (Hayes and Li, 2025). The modeled abundances indicate that the RT more closely reproduces the laboratory abundances, with deviations within $\pm 5\%$ for both H₂O ice and H₂SO₄·8H₂O with $\sim 100\ \mu\text{m}$ grains. In contrast, the LM shows slightly larger discrepancies, typically ranging from ± 5 – 15% from the true abundances. Interestingly, both LM and RT tend to consistently overestimate the abundance of H₂SO₄·8H₂O and underestimate H₂O ice across all mixtures. Nonetheless, when H₂SO₄·8H₂O either dominates ($> 80\%$ as observed on Europa’s trailing hemisphere; Carlson et al., 2005) or is present only in trace amounts ($\sim 10\%$ on areas in Europa’s leading hemisphere; Dalton III et al., 2013; Ligier et al., 2016), both the LM and RT render acceptable results within $\pm 10\%$ uncertainty. Thus, spectral modeling using the RT is preferred for constraining the surface composition across Europa, although the LM remains viable in specific compositional regimes.

Keywords: Spectral modeling, Radiative transfer, Europa, Icy moon, Surface composition, Planetary science

1. Introduction

Estimations of planetary surface composition—on both terrestrial and icy bodies—across the Solar System are done using spectral deconvolution via linear mixture (LM) modeling and radiative transfer (RT) theory, known as intimate mixture modeling (e.g., Li and Li, 2011; Dalton III et al., 2012; Li and Milliken, 2015; Goudge et al., 2015; Shirley et al., 2016; Liu et al., 2016; Emran et al., 2021; Emran and Stack, 2025; Emran, 2025a). In linear mixture modeling, also commonly known as the areal mixture approach, the reflectance of a mixture is represented as a linear combination of the spectra of its

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constituent materials, assuming that each material exists in spatially distinct patches (e.g., [Hapke, 2012](#); [Dalton III et al., 2012](#); [Emran et al., 2023](#); [Stack and Milliken, 2015](#)). In contrast, radiative transfer-based intimate mixture modeling assumes that the component materials are physically mixed in a “salt-and-pepper” fashion, resulting in a nonlinear combination of the reflectance spectra of the constituents (e.g., [Mustard and Pieters, 1987](#); [Poulet and Erard, 2004](#); [Li and Li, 2011](#); [Hapke, 2012](#); [Li and Milliken, 2015](#)). In planetary science, Hapke’s radiative transfer model [Hapke \(1981, 2001, 2012\)](#) has been extensively applied for spectral modeling to constrain the compositions of icy bodies in the outer Solar System (refer to [Emran and Chevrier, 2022, 2023](#); [Khuller and Emran, 2025](#), for exhaustive references).

Although both the LM and RT spectral modeling have their inherent advantages and limitations (e.g., [Keshava and Mustard, 2002](#); [Hapke, 2012](#)), they have been successfully applied to constrain the surface composition of Europa (e.g., [Shirley et al., 2016](#)). Using data from both ground- and space-based observations, Europa’s surface composition has been derived from visible and near-infrared (VIS-NIR) wavelengths using the LM modeling (e.g., [Dalton III, 2007](#); [Dalton III et al., 2012](#); [Shirley et al., 2010](#); [Brown and Hand, 2013](#); [Ligier et al., 2016](#); [Davis and Brown, 2024](#); [Emran, 2025a](#)) and the RT modeling based on [Hapke \(1981\)](#) theory (e.g., [Carlson et al., 2005](#); [Mishra et al., 2021](#); [Mermy et al., 2023](#)). While a variety of ice and non-ice components have been hypothesized on Europa’s surface, water (H_2O) ice and sulfuric acid octahydrate (SAO; $\text{H}_2\text{SO}_4 \cdot 8\text{H}_2\text{O}$) have consistently been identified as the major constituents using both the LM and RT modeling approaches, as well as by comparisons with laboratory spectra (e.g., [Moroz, 1966](#); [Pilcher et al., 1972](#); [Carlson et al., 1999a, 2005](#); [Dalton III, 2007](#); [Dalton III et al., 2012](#); [Shirley et al., 2010, 2016](#); [Ligier et al., 2016](#); [Mermy et al., 2023](#); [Emran, 2025a](#)).

Nonetheless, evaluating the accuracy of the LM and RT modeling in constraining surface composition remains a critical area of research for accurately understanding Europa’s surface and interpreting its geology. [Shirley et al. \(2016\)](#) compared abundance estimates derived from the LM and RT modeling using data from the Galileo Near-Infrared Mapping Spectrometer (NIMS; [Carlson et al., 1992](#)). Their analysis of spectra from various regions on Europa revealed that both approaches rendered consistent results, with linear correlation coefficients of ≥ 0.98 for H_2O ice and $\text{H}_2\text{SO}_4 \cdot 8\text{H}_2\text{O}$ abundances by both models ([Shirley et al., 2016](#)). However, until this current study, validation of these mixture modeling using laboratory data has remained lacking. Recently, [Hayes and Li \(2025\)](#) measured reflectance spectra of H_2O ice and $\text{H}_2\text{SO}_4 \cdot 8\text{H}_2\text{O}$ mixtures at varying ratios relevant to Europa—offering a unique opportunity to assess spectral modeling accuracy. In this study, as part of a series of investigations ([Emran, 2025b](#)), I assess the accuracy of the LM and RT modeling against the laboratory spectra of H_2O ice and $\text{H}_2\text{SO}_4 \cdot 8\text{H}_2\text{O}$ mixtures at near-infrared (NIR) wavelengths—providing a framework for quantitatively constraining Europa’s surface composition.

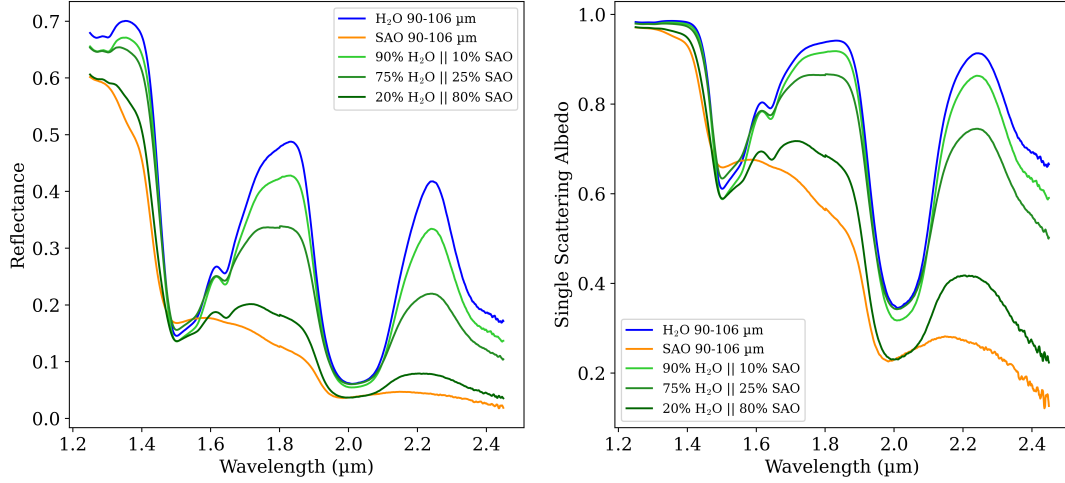


Figure 1: *Left panel:* Laboratory reflectance spectra of H_2O ice, $\text{H}_2\text{SO}_4 \cdot 8\text{H}_2\text{O}$ (sulfuric acid octahydrate; SAO), and their mixtures at varying proportions (%wt) and grain sizes of 90–106 μm (Hayes and Li, 2025). The blue spectrum represents the H_2O ice endmember, the orange spectrum represents the $\text{H}_2\text{SO}_4 \cdot 8\text{H}_2\text{O}$ endmember, and the spectra shown in shades of green correspond to mixtures of H_2O ice and $\text{H}_2\text{SO}_4 \cdot 8\text{H}_2\text{O}$ at different ratios. All reflectance spectra were collected from Hayes and Li (2025). *Right panel:* Single scattering albedo spectra corresponding to the reflectance spectra, derived using radiative transfer theory based on the Hapke (1981) model (refer to Section 2.2 for conversion details).

2. Observation and Methods

2.1. Laboratory spectra

Reflectance spectra of H_2O ice and $\text{H}_2\text{SO}_4 \cdot 8\text{H}_2\text{O}$ mixtures were measured at varying ratios and in the grain sizes of 90–106 μm (Hayes and Li, 2025). The measurements were conducted at ~ 77 K under standard viewing geometry, with an incidence angle (i) of 30° , an emission angle (e) of 0° , and a phase angle (g) of 30° at the spectral wavelength of ~ 1.25 – 2.45 μm . Although Europa’s surface temperature is ~ 120 K (e.g., Spencer et al., 1999; Carlson et al., 2009; Ashkenazy, 2019), reflectance spectra of H_2O ice and $\text{H}_2\text{SO}_4 \cdot 8\text{H}_2\text{O}$ measured at 77 K can be used to represent Europa’s surface relevant spectra, as spectral variations (absorption depths and position) over this temperature range are minor (e.g., Carlson et al., 2005; Mastrapa et al., 2008, 2009; Stephan et al., 2021; Hayes and Li, 2025). For details on the experimental setup and procedures, I refer readers to Hayes and Li (2025). The spectra of pure H_2O ice and pure $\text{H}_2\text{SO}_4 \cdot 8\text{H}_2\text{O}$ obtained by Hayes and Li (2025) are used as endmembers for both the LM and RT spectral modeling. The mixture dataset includes three compositions: 90% H_2O + 10% SAO, 75% H_2O + 25% SAO, and 20% H_2O + 80% SAO mixtures (by mass or %wt) at grain sizes of ~ 100 (90–106) μm (Fig. 1). The $\text{H}_2\text{SO}_4 \cdot 8\text{H}_2\text{O}$ grains used in this study are spherical, whereas the H_2O ice grains are angular to sub-angular (T. W. Hayes, 2025; personal communication).

2.2. LM and RT modeling

In linear mixture modeling, the mixture spectrum is represented as the weighted sum of the reflectances of each endmember, with the weights corresponding to their abundances (e.g., [Emran et al., 2023](#)):

$$r_m = \sum_{i=1}^k f_i \cdot r_i ; 0 \leq f_i \leq 1 ; \sum_{i=1}^k f_i = 1 \quad (1)$$

where r_m is the reflectance of the mixture spectrum, r_i is the reflectance of the i^{th} constituent endmember, and f_i is the fractional abundance (%wt) of the i^{th} endmember (e.g., [Dalton III et al., 2012](#); [Emran et al., 2021](#); [Emran and Stack, 2025](#)). Using Eq. (1), I estimated the modeled abundances of H_2O ice and $\text{H}_2\text{SO}_4 \cdot 8\text{H}_2\text{O}$ for reflectance spectra of mixtures.

In contrast, intimate mixture modeling based on radiative transfer accounts for the non-linear scattering properties of reflectance, but assumes linear behavior when reflectance is converted to single scattering albedo ([Hapke, 1981, 2001, 2012](#); [Mustard and Pieters, 1989](#); [Mustard and Glotch, 2019](#)). Therefore, I converted both the endmember and mixture reflectance spectra to single scattering albedo using the Hapke theory ([Hapke, 1981](#)) with the same viewing geometry and grain size as the measured reflectance spectra (Fig. 1). The reflectance (also known as radiance coefficient) in the Hapke model ([Hapke, 1981](#)) can be expressed as:

$$r(i, e, g) = \frac{w}{4} \cdot \frac{1}{\mu_e + \mu_i} \{ [1 + B(g)] P(g) + H(\mu_e)H(\mu_i) - 1 \} \quad (2)$$

where r is the reflectance, w is the average single-scattering albedo, μ_i and μ_e are the cosines of incidence angle (i) and emission angle (e), respectively, $B(g)$ is the backscattering function, $P(g)$ is the single-particle phase function, and $H(\mu_x)$ is the multiple scattering H functions of [Chandrasekhar \(1960\)](#). Refer to relevant literature for details of these functions (e.g., [Hapke, 1981, 2001, 2012](#)). However, for the laboratory spectra considered in this study, I assume no backscattering ($B(g) = 0$) and an isotropic phase function ($P(g) = 1$), following [Lapotre et al. \(2017\)](#). Under these assumptions, the Hapke reflectance equation simplifies to (e.g., [Mustard and Pieters, 1987, 1989](#)):

$$r(i, e, g) \approx \frac{w}{4} \cdot \frac{1}{\mu_e + \mu_i} \cdot H(\mu_e)H(\mu_i) \quad (3)$$

The validity of Eq. 3 for abundance estimation has been tested against laboratory reflectance spectra of intimate mineral mixtures, under the assumption of isotropic scattering behavior of particulate media at phase angles of $15^\circ - 40^\circ$ ([Mustard and Pieters, 1987, 1989](#)). At these intermediate phase angles, the backscattering is assumed negligible ($B(g) \ll 1$) and becomes significant only at very small phase angles ($< 10^\circ$; [Mustard and Pieters 1987](#)). Using this simplified method (Eq. 3), abundance estimates are accurate to within a few percent (%) when particle grain sizes are known ([Mustard and Pieters, 1987, 1989](#)).

The Chandrasekhar H functions describing multiple scattering can be derived as (Hapke, 2002):

$$H(x) \approx \left[1 - wx \left(r_o + \frac{1 - 2r_o x}{2} \ln \left(\frac{1 + x}{x} \right) \right)^{-1} \right] \quad (4)$$

$$r_o = \frac{1 - \gamma}{1 + \gamma} \quad (5)$$

$$\gamma = \sqrt{1 - w} \quad (6)$$

where r_o is the bihemispherical reflectance for isotropic scatterers, γ is the albedo factor, and x represents either the cosine of the incidence angle ($\mu_i = \cos(i)$) or the emission angle ($\mu_e = \cos(e)$).

The photon penetration depth (δ) into a particulate surface or media is inversely proportional to the imaginary part (k) of the material's optical constant at a given wavelength (e.g., Born et al., 1999; Hapke, 2012). For both H_2O ice and $\text{H}_2\text{SO}_4 \cdot 8\text{H}_2\text{O}$, the low k values at NIR (e.g., Mastrapa et al., 2008, 2009; Carlson et al., 2005) result in δ on the order of hundreds of micrometers. Because the grain diameter ($d \sim 100\mu\text{m}$) is much smaller than the penetration depth ($d < \delta$), photons interact with multiple grains, making multiple scattering the dominant effect (e.g., Hapke, 1981, 2012). This process can induce significant spectral changes in intimate mixture even with trace amounts of material (e.g., Hayes and Li, 2025). Thus, multiple scattering terms (H -functions) were included in RT modeling (Eq. 3) to ensure accurate abundance estimates.

Upon conversion to single scattering albedo from the laboratory reflectance spectra using Eqs. (3)–(6), the RT spectral modeling can be expressed as follows (e.g., Stack and Milliken, 2015; Goudge et al., 2015; Lapotre et al., 2017):

$$w_m = \sum_{i=1}^k f_i \cdot w_i \quad (7)$$

where w_m is the single scattering albedo of the mixture spectrum, w_i is the single scattering albedo of the i^{th} constituent endmember, and f_i is the relative fractional (geometric) cross section of the i^{th} endmember. The relative geometric cross-section (also known as F-parameter) is a function of the mass fraction, density, and particle diameter of components in a mixture (e.g., Mustard and Pieters, 1987, 1989).

While grain size has a major influence on the spectral characteristics of materials, grain shape can also affect the resulting spectra (e.g., Shkuratov and Grynko, 2005; Hapke, 2012). However, since both endmembers— H_2O ice and $\text{H}_2\text{SO}_4 \cdot 8\text{H}_2\text{O}$ —were modeled with the same grain size ($\sim 100\mu\text{m}$), shape effects were ignored. Under simplified assumptions, such as similar grain sizes and densities of particles, f_i can be considered equivalent to the fractional contribution or abundance (%wt) of the corresponding endmember (e.g., Mustard and Pieters, 1987). However, the densities of H_2O ice and SAO differ slightly under outer Solar System-relevant temperatures— H_2O ice has a density of $\sim 1 \text{ g/cm}^3$ (e.g., Blackman and Lisgarten, 1957; Yarnall and Hudson, 2022), whereas the

slightly denser $\text{H}_2\text{SO}_4 \cdot 8\text{H}_2\text{O}$ has a density of $\sim 1.3 \text{ g/cm}^3$ (Maynard-Casely et al., 2012, 2013). Accordingly, the relative fractional cross-section for H_2O ice and $\text{H}_2\text{SO}_4 \cdot 8\text{H}_2\text{O}$ was scaled by their respective density factors in Eq. (7), and the fractional abundances in the mixtures were adjusted accordingly.

In this study, I employed the Markov Chain Monte Carlo (MCMC) simulation technique (Hogg and Foreman-Mackey, 2018) to estimate the model abundances using Equations (1) and (7), following the approach outlined in Emran et al. (2021). The MCMC technique used here adopts a Bayesian inference framework to predict the model parameters, e.g., the abundances of each endmember. Using the *emcee* Python package (Foreman-Mackey et al., 2013), I implemented the models and performed 1,000 iterations to estimate the parameters.

3. Results

Fig. 2 presents the modeled spectra generated using the linear mixture (left panel) and radiative transfer (right panel) modeling approaches, along with the estimated best-fit mean abundances (Est.) for H_2O ice and $\text{H}_2\text{SO}_4 \cdot 8\text{H}_2\text{O}$ across different laboratory mixtures. The estimated abundances vary with mixture ratios (%wt), as shown in the row-wise subplots (Fig. 2).

For the 90% H_2O || 10% SAO mixture (top row in Fig. 2), the RT modeling results in better abundance estimates of H_2O ice and $\text{H}_2\text{SO}_4 \cdot 8\text{H}_2\text{O}$ than the LM modeling, even though the LM modeling produces a lower root mean square error (RMSE) value—0.008 (LM) vs 0.011 (RT). Specifically, the LM and RT estimated mean $\pm 1\sigma$ abundances are $82 \pm 4\%$ H_2O || $18 \pm 4\%$ SAO and $89 \pm 3\%$ H_2O || $11 \pm 3\%$ SAO, respectively (Table 1). Thus, the mean abundance offset from the true composition is $\pm 8\%$ for LM and $\pm 1\%$ for RT modeling (Fig. 3).

In the case of the 75% H_2O || 25% SAO mixture (middle row in Fig. 2), the abundance estimates by the RT modeling again outperform the LM modeling results. Moreover, the RT modeling also shows a lower RMSE value than the LM modeling—0.011 vs 0.018, respectively. The estimated abundances are $60 \pm 4\%$ H_2O || $40 \pm 4\%$ SAO for LM and $72 \pm 4\%$ H_2O || $28 \pm 4\%$ SAO for RT modeling (Table 1), resulting in mean abundance offsets of $\pm 15\%$ (LM) and $\pm 3\%$ (RT) from the true values (Fig. 3).

Lastly, for the 20% H_2O || 80% SAO mixture (bottom row in Fig. 2), both models perform comparably in terms of RMSE value (0.012 by RT and 0.011 by LM), though the RT modeling still results in closer abundance estimates to the true abundance. The LM and RT estimates are $14 \pm 4\%$ H_2O || $86 \pm 4\%$ SAO and $18 \pm 5\%$ H_2O || $82 \pm 5\%$ SAO, respectively (Table 1), corresponding to mean abundance offsets of $\pm 6\%$ and $\pm 2\%$, respectively (Fig. 3).

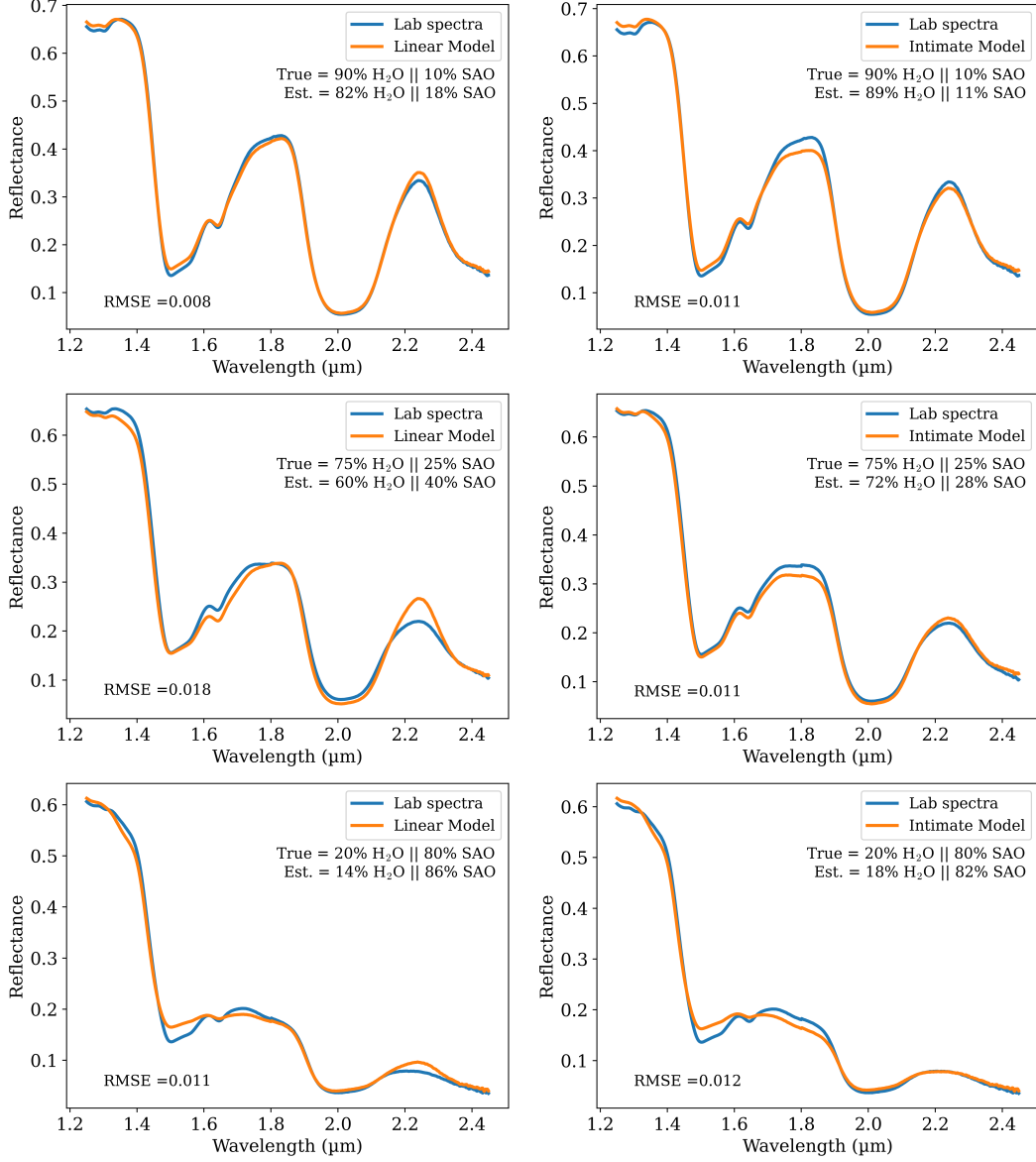


Figure 2: Comparison of linear mixture modeling (left panel) and intimate mixture modeling using the Hapke (1981) radiative transfer theory (right panel). In each subplot, the blue spectrum represents the laboratory reflectance spectrum of the H₂O–SAO mixture (Hayes and Li, 2025), while the orange spectrum shows the modeled spectrum from either the LM or Hapke-based RT modeling. The true (laboratory) and estimated (model-derived) mean abundances of H₂O ice and sulfuric acid octahydrate (SAO) are indicated in each subplot, along with the root mean square error (RMSE) of the spectral fit.

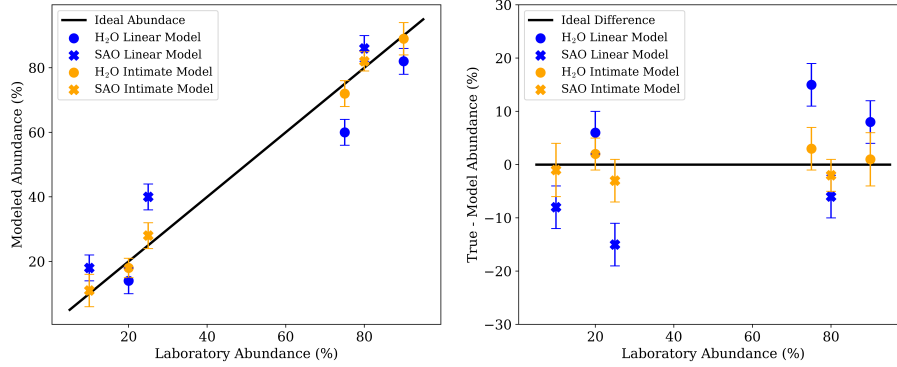


Figure 3: Comparison of modeled abundances (mean $\pm 1\sigma$) of H₂O ice and sulfuric acid octahydrate (SAO) using the linear mixture modeling and the intimate mixture modeling based on the Hapke (1981) radiative transfer theory. *Left panel*: Modeled abundances plotted against the expected (ideal) laboratory abundances. Blue and orange markers represent estimates from the LM and Hapke-based RT modeling, respectively. Circle and cross markers represent H₂O ice and SAO abundances, respectively, by the corresponding mixture modeling. *Right panel*: Difference between the laboratory (true) abundances and the modeled abundances (True-Model abundance) for each component using both modeling approaches (LM and RT).

Table 1: Modeled abundances of H₂O ice and sulfuric acid octahydrate (SAO) using linear mixture modeling and intimate mixture modeling based on the Hapke (1981) radiative transfer theory. Abundances are reported as fractional percentages (%wt) as mean $\pm 1\sigma$.

Mixture abundance (laboratory mixture)	Linear mixture (LM) abundance		Intimate mixture (RT) abundance	
	H ₂ O (%)	SAO (%)	H ₂ O (%)	SAO (%)
90% H ₂ O + 10% SAO	82 $\pm$ 4	18 $\pm$ 4	89 $\pm$ 3	11 $\pm$ 3
75% H ₂ O + 25% SAO	60 $\pm$ 4	40 $\pm$ 4	72 $\pm$ 4	28 $\pm$ 4
20% H ₂ O + 80% SAO	14 $\pm$ 4	86 $\pm$ 4	18 $\pm$ 5	82 $\pm$ 5

4. Discussion and implications for Europa

In this study, I compared two widely-used spectral modeling approaches— linear mixture modeling and radiative transfer-based intimate mixture modeling —against laboratory spectra of H₂O ice and H₂SO₄·8H₂O mixtures at varying ratios (%wt) relevant to Europa (Hayes and Li, 2025). The results suggest several key insights regarding the future implementation of spectral modeling (LM and RT) for constraining surface compositional abundances, and by extension, for interpreting the geological processes and surface compositional evolution on Jupiter’s moon.

This study demonstrates that the spectral fit quality—quantified by metrics such as root mean square error (RMSE) or chi-square (χ^2) values—does not necessarily reflect the accuracy of the estimated compositional abundances. For example, in the 90% H₂O || 10% SAO mixture case, the LM model resulted in a lower RMSE compared to the RT model (also, the LM visually shows a closer spectral fit to the laboratory spectra), yet

the RT approach produced significantly more accurate abundance estimates. Conversely, in another case (75% H₂O || 25% SAO mixture), a better spectral fit—as indicated by a lower RMSE—did correspond to abundance estimates that closely matched the true compositional abundance. This suggests that while a lower RMSE or χ^2 value may sometimes reflect improved abundance accuracy, it is not a consistent indicator or does not guarantee the model’s performance in compositional retrieval (abundance estimates of surface materials) on Europa.

Across all tested mixture ratios, the RT modeling consistently produces mean abundance estimates that are within $\pm 5\%$ of the true laboratory values for H₂O ice and SAO mixtures with $\sim 100\ \mu\text{m}$ grains. In contrast, the LM modeling tends to result in larger deviations, typically in the range of $\pm 5\text{--}15\%$. This observation may suggest that Europa’s surface is more accurately represented by an intimate mixture (RT) modeling, regardless of geologic units (e.g., Greeley et al., 2000; Prockter and Patterson, 2009; Daubar et al., 2024; Leonard et al., 2024) across the leading and trailing hemispheres, given that the grain sizes are comparable ($\sim 100\ \mu\text{m}$). However, the deviation in the LM estimates is relatively smaller when sulfuric acid abundance dominates in the mixture. For instance, the LM estimated abundances deviate by only $\pm 6\%$ from the true values in a high-SAO scenario (80% in mixture). A similar deviation in abundance ($\pm 8\%$) is also observed where sulfuric acid is present in trace amounts, such as in $\sim 10\%$ of the mixture. Thus, in regions where surface materials show a patchy distribution and the abundance of sulfuric acid is either very high ($> 80\%$) or very low ($\sim 10\%$), the LM model produces viable abundance estimates with $\sim 100\ \mu\text{m}$ grains. For instance, a surface composed of patches of dark and bright materials on Europa (Fig. 4)—arranged in a checkerboard-like distribution—may favor the application of the LM model (e.g., Shirley et al., 2016), as this approach is based on areal mixing of materials. In contrast, RT modeling is better suited for regions where surface materials are thoroughly mixed at the particle level—similar to a lunar regolith-type intimately mixed surface (e.g., Shkuratov et al., 1999)—rather than being segregated into distinct patches, which could result from thermal segregation of surface materials (Spencer, 1987; Sorli and Hayne, 2023). RT modeling may particularly be relevant in areas influenced by radiolytic processes on Europa (Fig. 4; Shirley et al., 2016), where radiation likely interacted uniformly with the molecules in the surface regolith of the moon (Carlson et al., 1999a, 2002). Nonetheless, regardless of the magnitude of deviation, both LM and RT modeling consistently overestimates the abundance of H₂SO₄·8H₂O and underestimating H₂O ice (Table 1). These discrepancies are consistently observed across all laboratory mixture cases analyzed in this study (see Fig. 3).

Photons at NIR wavelengths penetrate a finite depth into a particulate surface (e.g., Born et al., 1999). This penetration depth is a critical factor, as it governs multiple scattering effects in intimate mixtures or granular media (Hapke, 1981, 2012). The areal mixing approach assumes that the surface is composed of patches with uniform composition and does not account for variations in depth within grains or regolith. Under this assumption, each photon interacts with only a single endmember, thereby neglecting multiple scattering between different compositional units or grains. In contrast, intimate mixtures exhibit highly nonlinear spectral behavior at NIR wavelengths be-

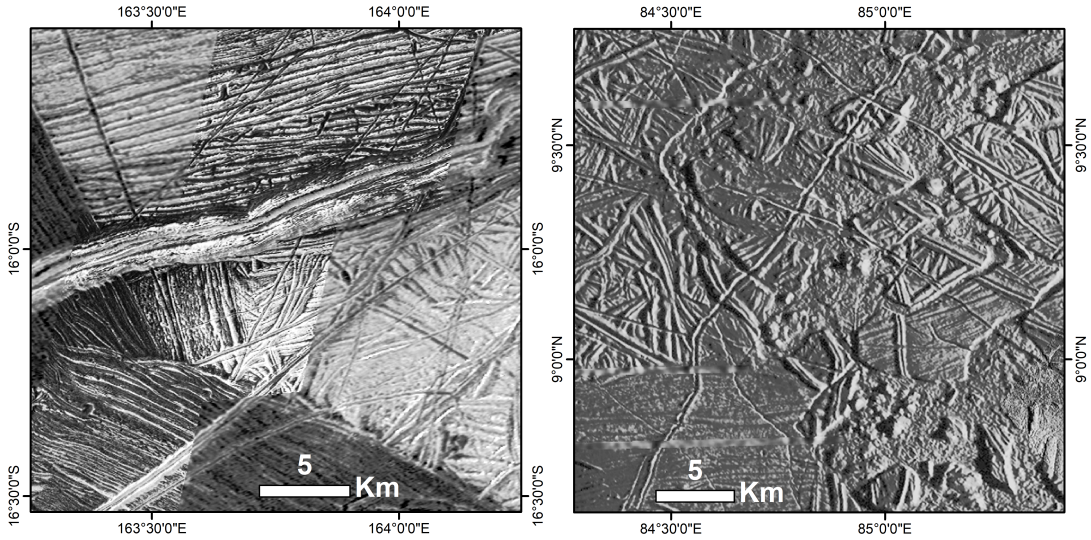


Figure 4: *Left panel:* Patches of dark and bright materials on Europa observed in a high-resolution image from the Galileo Solid State Imaging (SSI) system (Belton et al., 1992). The image is centered at approximately 16°S, 163°45'E, with a spatial resolution of ~ 50 m/pixel. *Right panel:* A high-resolution Galileo/SSI image of Europa's chaos terrains on the trailing hemisphere (near center)—an area exposed to radiolytic processes (e.g., Carlson et al., 1999a, 2002)—showing no discernible dark and bright patches. The image is centered at approximately 9°15'N, 84°50'E, with a spatial resolution of ~ 50 m/pixel. Both SSI images are adopted from Malaska et al. (2024).

cause light is multiply scattered among mineral species or grains (Hapke, 1981, 2012; Mustard and Pieters, 1987, 1989). As a result, even trace amounts of material can significantly influence the observed spectra of an intimate mixture (e.g., Hapke, 2012; Hayes and Li, 2025). These limitations of the LM modeling approach—which ignores photon penetration depth, multiple scattering, and compositional variation with depth—likely contribute to the higher discrepancies in abundance estimates of mixture components, particularly when grain sizes are much smaller than the photon penetration depth. By comparison, RT modeling accounts for photon penetration depth by incorporating multiple scattering effects when estimating the complex spectral properties of mixtures. While multiple scattering strongly influences reflectance spectra used in LM modeling, its effects are reduced when reflectance is converted to single scattering albedo (e.g., Hapke, 1981, 2012). In this form, the single scattering albedo spectra of a mixture is approximately linearly proportional to the relative fractional abundances of the constituent materials (e.g., Hapke, 1981, 2012; Mustard and Glotch, 2019). Thus, by properly accounting for the complexity of light scattering in granular or particulate media, RT modeling provides more accurate abundance estimates—within $\pm 5\%$ in this study—compared to LM modeling.

The performance of the LM and RT modeling was compared here using laboratory spectra of H_2O ice and sulfuric acid with grain sizes ranging from 90–106 μm (Hayes and Li, 2025). Due to the limited availability of laboratory spectra for other grain sizes for

H₂O and H₂SO₄·8H₂O mixtures, the influence of grain size on the modeling performance for these mixtures could not be directly assessed. However, the effect of grain size on reflectance spectra—and thus on spectral modeling—cannot be overlooked, as it is well-established that grain size significantly affects spectral features (e.g., [Stephan et al., 2021](#); [Mastrapa et al., 2008](#); [Emran and Chevrier, 2023](#)). Relatively small grain sizes—ranging from 20–30 μm on the leading hemisphere to ~ 100 μm near the poles of the trailing hemisphere—have been suggested to characterize Europa’s surface (e.g., [Kieffer and Smythe, 1974](#); [Hansen and McCord, 2004](#); [Carlson et al., 2005, 2009](#); [Moore et al., 2009](#); [Becker et al., 2024](#)). However, spectral analyses suggest the presence of coarser grains—up to ~ 1 mm in size—particularly in regions such as large lineae, microchaos, and band geologic units ([Ligier et al., 2016](#); [Emran, 2025a](#)). Thus, the findings presented here are most applicable to regions on Europa where grain sizes are relatively small, around ~ 100 μm . In contrast, terrains associated with effusive cryovolcanism (viz. surface flow) or surface–subsurface exchange, such as microchaos, are suggested to host coarser grains, potentially up to ~ 1 mm in size ([Ligier et al., 2016](#); [Emran, 2025a](#)). Consequently, I acknowledge that the results of this study have not been validated for those regions with coarser surface materials. A separate investigation using laboratory spectra with a broader range of grain sizes—including mixtures containing both micron- and millimeter-scale diameter grains—is warranted to assess the impact of grain size on spectral modeling accuracy. Notably, a detailed study comparing the accuracy of LM and RT modeling for H₂O ice mixtures with coarse (~ 1 mm) and fine (~ 100 μm) grains relevant to Europa is presented in [Emran \(2025b\)](#).

Moreover, this study investigated the accuracy of LM and RT modeling using the reflectance spectra of the binary mixtures of H₂O ice and SAO—the two major components identified on Europa (e.g., [Moroz, 1966](#); [Pilcher et al., 1972](#); [Carlson et al., 1999a, 2005, 2009](#); [Becker et al., 2024](#))—at different mixing ratios ([Hayes and Li, 2025](#)). However, Europa’s surface is also hypothesized to host other species, such as CO₂ ([Villanueva et al., 2023](#)), H₂O₂ ([Carlson et al., 1999b](#)), SO₂ ([Lane et al., 1981](#)), NH₃-bearing compounds ([Emran, 2025a](#)), and various hydrated sulfate and chlorine (Cl-bearing) salts (e.g., [Davis and Brown 2024](#); [Becker et al. 2024](#), and references therein). To comprehensively assess the performance of LM and RT modeling, it is important to validate these approaches against mixtures that include a broader suite of materials. While laboratory spectra of these species would be invaluable for such testing, well-constrained reflectance spectra of intimate mixtures are limited and currently available for binary mixtures of H₂O ice and H₂SO₄·8H₂O. Thus, I emphasize the need for a robust set of optical constants and laboratory mixture spectra for these additional constituents, as they are essential for a complete understanding of Europa’s surface composition as constrained through spectral modeling using either areal mixing or radiative transfer approaches.

The use of zero backscattering and isotropic scattering in this study was based on the controlled viewing geometry (phase angle of 30°) of the laboratory spectra and the known grain sizes of the materials. For remotely sensed data, such as from telescopes or spacecraft, however, the scenario differs—viewing geometry varies from observation to observation, and the grain sizes of the regolith are often poorly constrained or un-

known. In such cases, photometric analyses of planetary surfaces are performed using the Hapke radiative transfer model (Hapke, 1981) to estimate surface backscattering and phase parameters, as done for Europa (e.g., Buratti and Veverka, 1985; Buratti, 1995; Domingue and Hapke, 1992; Belgacem et al., 2018, 2020). Accurate backscattering and phase parameters are therefore essential for constraining reliable abundance estimates of surface compositions using RT modeling from ground- and space-based observations.

Nonetheless, both LM and RT models produce acceptable results— within $\pm 10\%$ uncertainty— under specific compositional conditions (H_2O and $\text{H}_2\text{SO}_4 \cdot 8\text{H}_2\text{O}$ mixtures with grain sizes around $\sim 100 \mu\text{m}$) relevant to regions on Europa. This includes the regions where $\text{H}_2\text{SO}_4 \cdot 8\text{H}_2\text{O}$ is the dominant component ($> 80\%$ abundance), such as the trailing hemisphere (e.g., Carlson et al., 2005), as well as regions where $\text{H}_2\text{SO}_4 \cdot 8\text{H}_2\text{O}$ occurs in trace amounts ($\sim 10\%$ abundance), as observed in some parts of the leading hemisphere (Dalton III et al., 2013; Ligier et al., 2016). Overall, the results support the use of radiative transfer-based intimate mixture spectral modeling (e.g., Hapke, 1981; Shkuratov et al., 1999) as the preferred method for constraining the surface composition of Europa. However, linear mixture modeling remains a viable approach in compositional regimes where the abundance of sulfuric acid is either very high or very low, allowing for flexibility depending on the available data and modeling constraints. These findings provide a valuable guideline for future compositional analyses using imaging spectrometers, such as ESA-JUICE’s Moons and Jupiter Imaging Spectrometer (MAJIS; Poulet et al., 2024) and NASA-Europa Clipper’s Mapping Imaging Spectrometer for Europa (MISE; Blaney et al., 2024), for an improved understanding of the composition of Europa’s surface.

Notations

r = Reflectance or radiance factor
 r_m = Reflectance of mixture
 r_o = Bihemispherical reflectance for isotropic scatterers
 f = Fractional abundance/ factor
 w = Single scattering albedo (SSA)
 w_m = SSA of the mixture spectrum
 i = Incidence angle
 e = Emission angle
 g = Phase angle
 $B(g)$ = Backscattering function
 $P(g)$ = Single-particle phase function
 $H(x)$ = Chandrasekhar's H functions
 μ_i = Cosine of the incidence angle
 μ_e = Cosine of the emission angle
 γ = Albedo factor
 δ = Photon penetration depth
 k = Imaginary part of optical constants
 d = Diameter or Grain Size

Data Availability

The laboratory spectral data used in this study were collected from [Hayes and Li \(2025\)](#) and can be accessed at <https://doi.org/10.5281/zenodo.13852483>. Galileo Solid State Imaging (SSI) data used in this study can be found in the National Aeronautics and Space Administration's Planetary Data System: Imaging Node Server at <https://pdsimage2.wr.usgs.gov/> ([Malaska et al., 2024](#)).

Declaration of generative AI

During the preparation of this work the author(s) used OpenAI in order to improve the readability and language of the manuscript. After using this tool/service, the author(s) reviewed and edited the content as needed and take(s) full responsibility for the content of the publication.

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